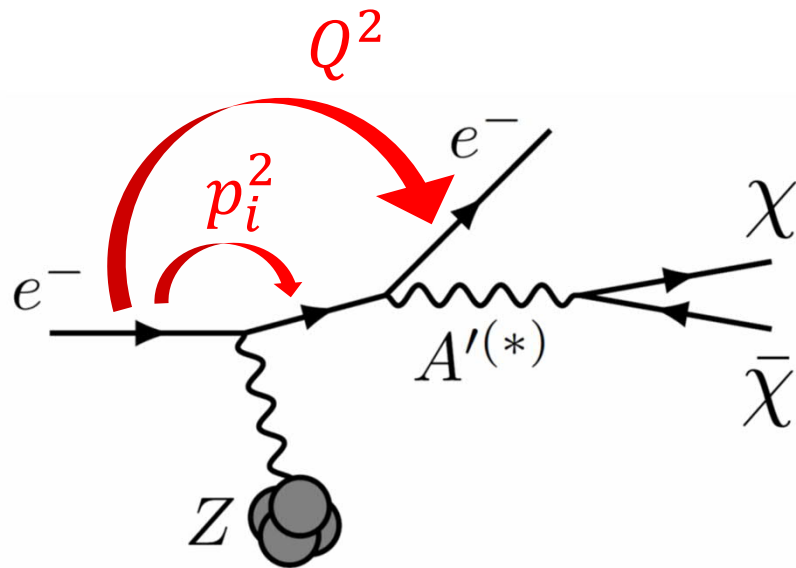


LDM events – understanding backgrounds



Momentum transfers

p_i^2 *internal scale for target interaction in signal event*

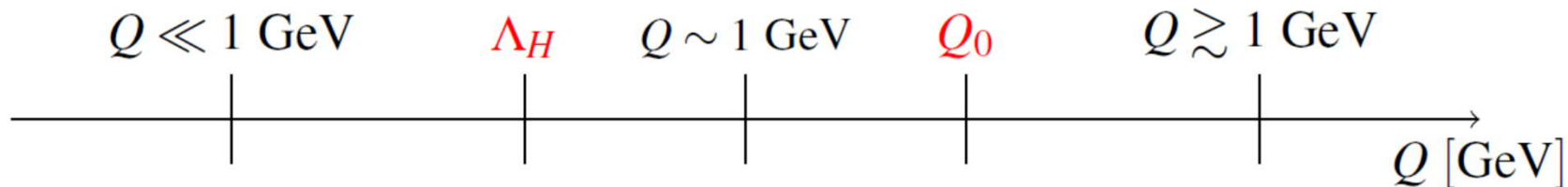
Q^2 *measured main scale in background events*



Momentum transfer scale



Quantum states & theory



hadrons
ChPT

hadrons → partons
form factors + starting PDFs

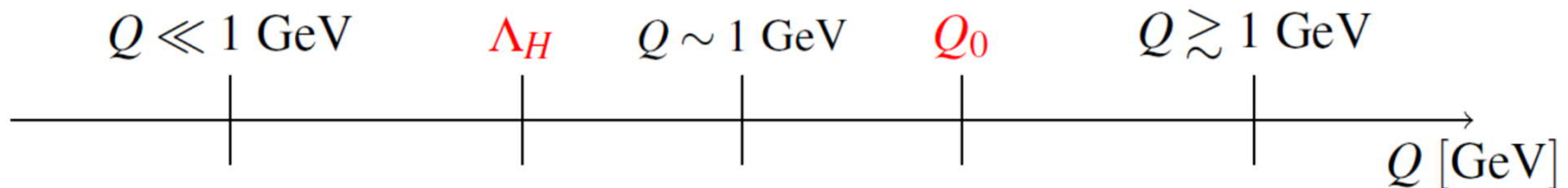
partons
pQCD

Understanding PDFs and the proton spin puzzle

A. Ekstedt, H. Ghaderi, G. Ingelman, S. Leupold

arXiv:1807.06589 Nucleon parton distributions from hadronic quantum fluctuations

arXiv:1808.06631 Towards solving the proton spin puzzle



hadronic fluctuations

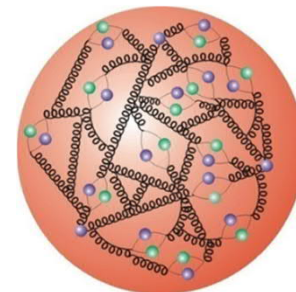
model for PDFs at Q_0

partonic fluctuations $|uud q\bar{q}\rangle$

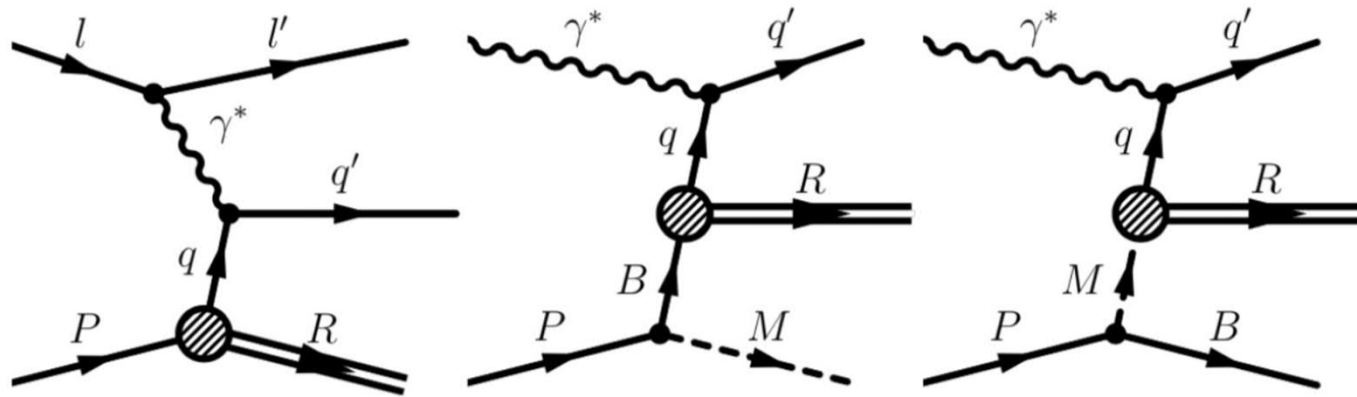
DGLAP Q^2 evolution of PDFs

$$|p\rangle = \alpha_0 |p_0\rangle + \alpha_{p\pi} |p\pi^0\rangle + \alpha_{n\pi} |n\pi^+\rangle + \dots$$

p-wave \rightarrow L=1 not observed in DIS



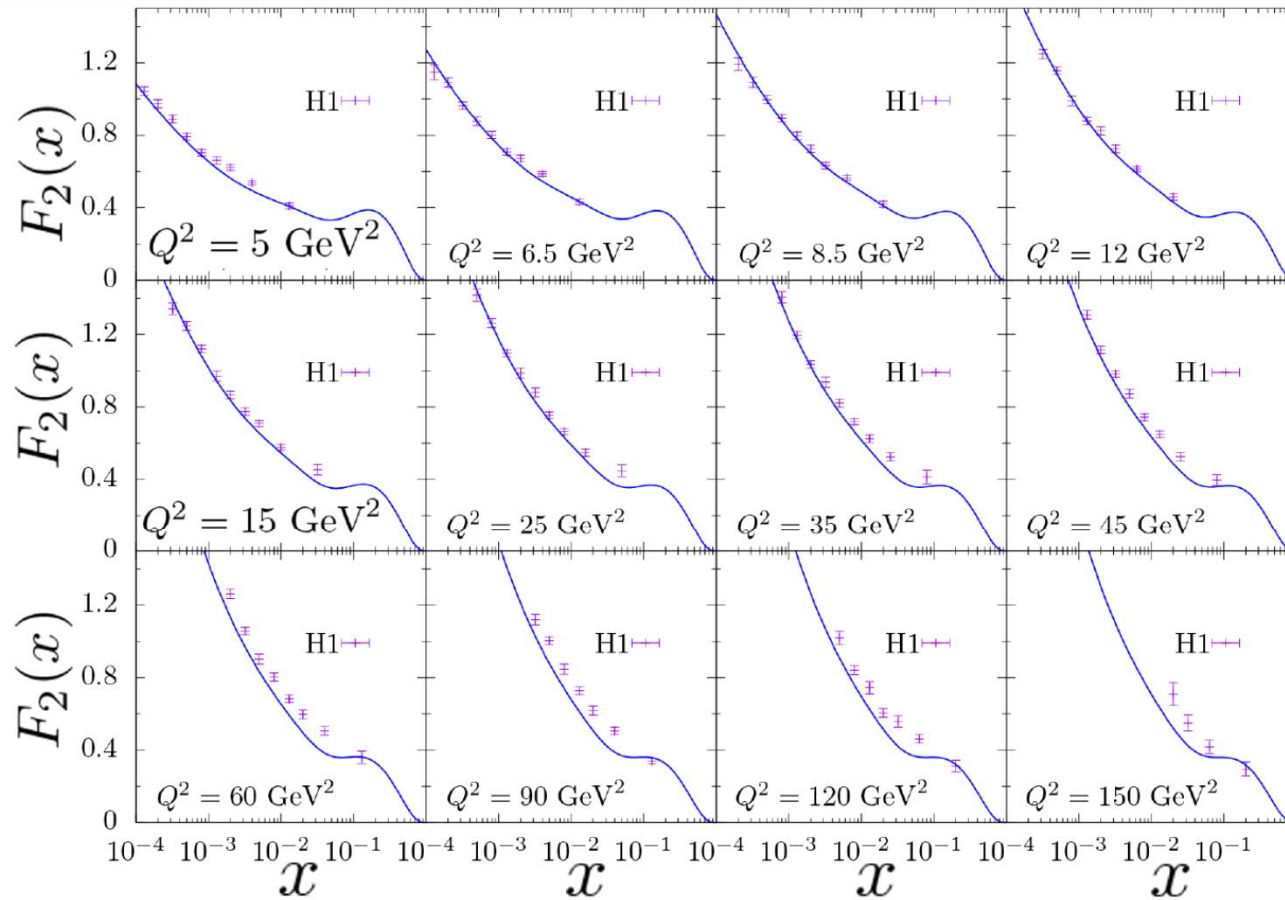
Fock expansion of proton state $|P\rangle = \sqrt{Z} |P\rangle_{\text{bare}} + \alpha_{n\pi^+} |n\pi^+\rangle + \alpha_{P\pi^0} |P\pi^0\rangle$
 $+ \alpha_{\Delta^{++}\pi^-} |\Delta^{++}\pi^-\rangle + \alpha_{\Delta^+\pi^0} |\Delta^+\pi^0\rangle + \alpha_{\Delta^0\pi^+} |\Delta^0\pi^+\rangle$
 $+ \alpha_{\Lambda K^+} |\Lambda K^+\rangle + \dots$



$$f_{i/P}(x) = f_{i/P}^{\text{bare}}(x) + \sum_{\lambda, H \in \{B, M\}} \int dy dz \delta(x - yz) f_{i/H}^{\text{bare}}(z) f_{H/P}^{\lambda}(y)$$

- LO ChPT Lagrangian $\rightsquigarrow f_{H/P}^{\lambda}(y) = \frac{|g_{BM}|^2}{(2\pi)^3 2y(1-y)} \int d^2 k_{\perp} \left| \phi \frac{S^{\lambda}(y, \mathbf{k}_{\perp})}{m_p^2 - m^2(y, k_{\perp}^2)} \right|^2$
 $y = p_B^+ / p_P^+ \quad \phi(y, k_{\perp}^2, \Lambda_H^2) = \exp[-(\mathbf{p}_M^2 + \mathbf{p}_B^2) / (2\Lambda_H^2)]$
- $f_{i/H}^{\text{bare}}(x) = \int \frac{d^4 k}{(2\pi)^4} \delta(k^+ / p_H^+ - x) \exp\left[-\frac{(k_0 - m_i)^2 + k_x^2 + k_y^2 + k_z^2}{2\sigma_i^2}\right]$ at $Q = Q_0$
DGLAP $Q > Q_0$
- 'Free' parameters: $\sigma_1, \sigma_2, \sigma_g$ and Λ_H, Q_0

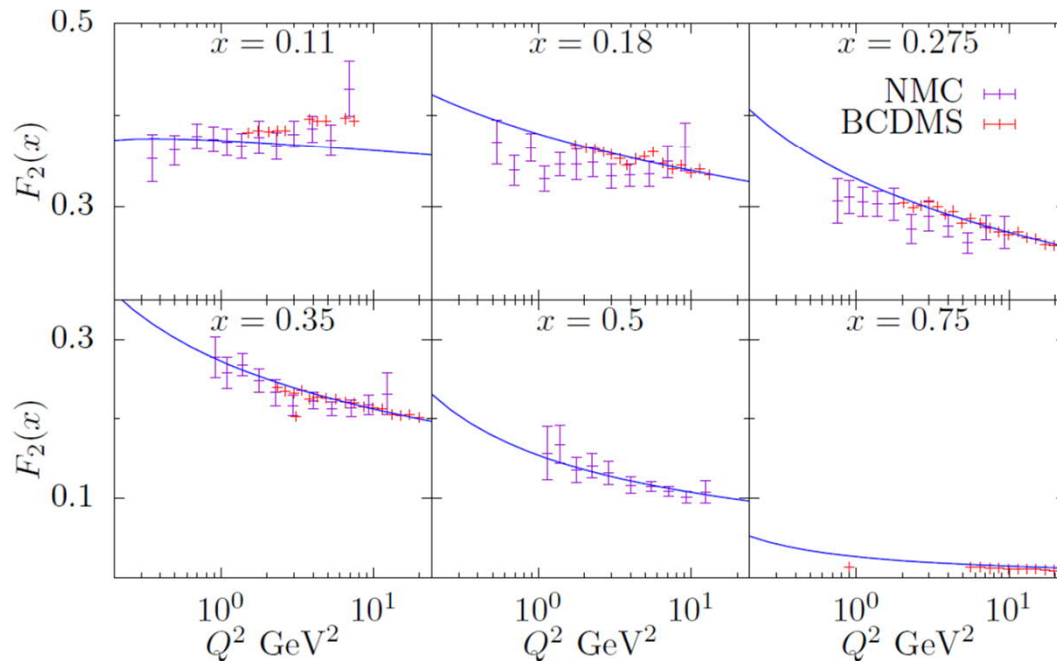
A priori expected value: $\sigma_i \sim 1/D_H \sim 0.1$ GeV and $\Lambda_H \sim Q_0 \sim 0.5$ GeV



Unpolarised
structure functions
↓
parameter values

F_2 at small x
→ σ_g, Q_0

F_2, F_3 at larger x
→ $\sigma_1, \sigma_2, \Lambda_H$



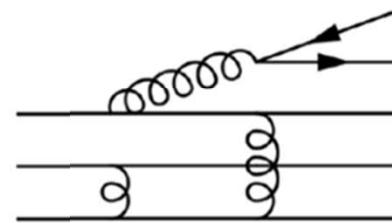
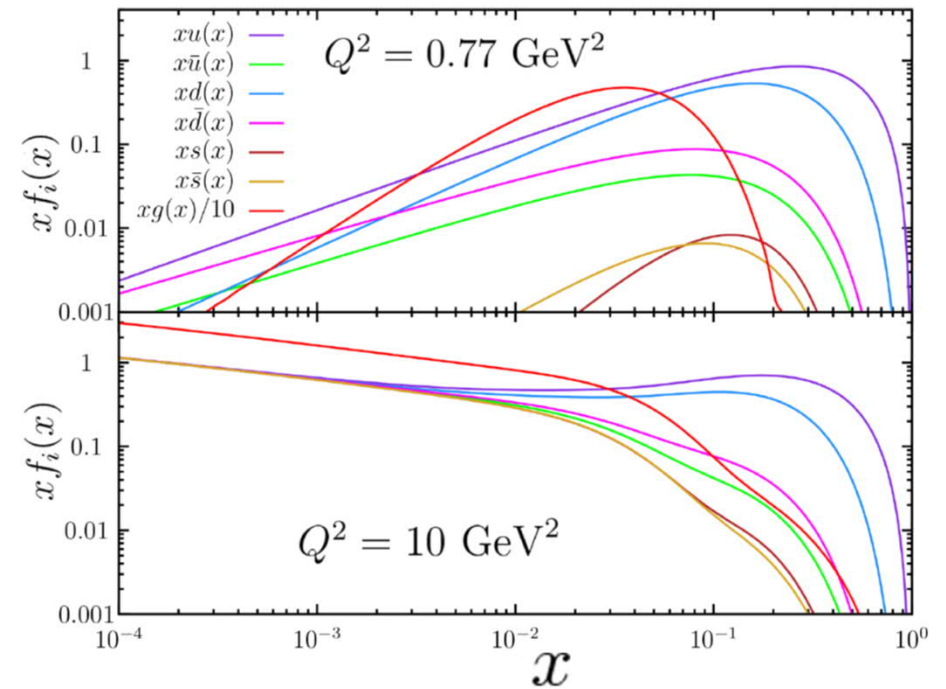
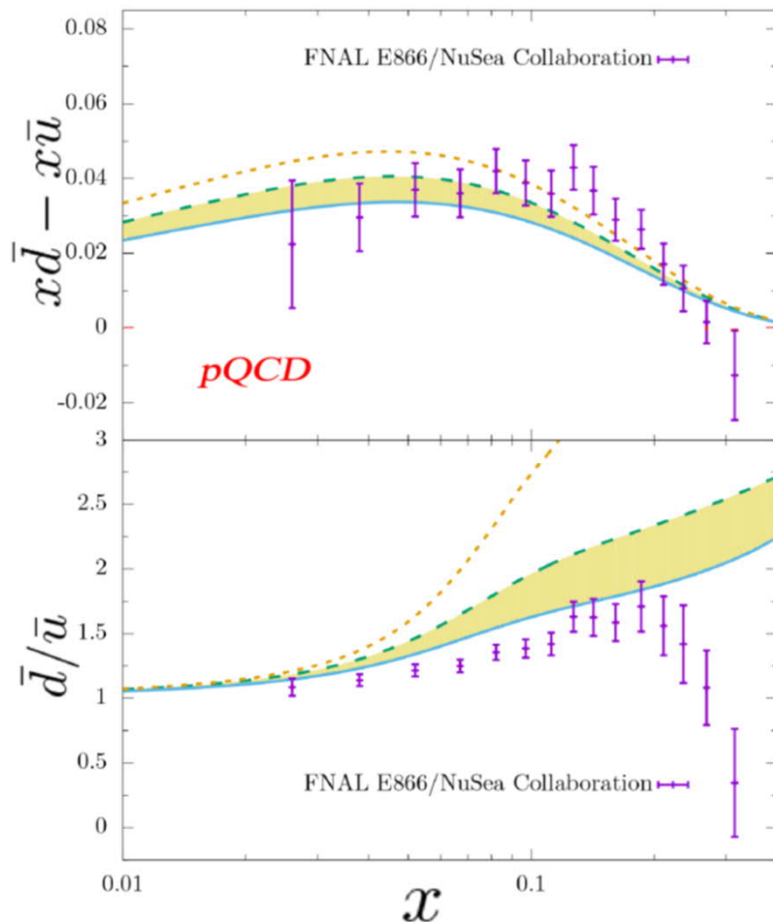
$\sigma_g = 0.028 \text{ GeV},$!! ??
 $\sigma_1 = 0.11 \text{ GeV},$
 $\sigma_2 = 0.22 \text{ GeV},$ Pauli principle !?

$\Lambda_H = 0.87 \text{ GeV},$ hadron-parton
 $Q_0 = 0.88 \text{ GeV}$ transition !!

Resulting unpolarised PDFs
based on 5 physical parameters
→ understanding!

cf. conventional PDF fits with
~30 parameters

Also explains sea quark asymmetries:



- From pQCD point of view: since $m_u, m_d \ll \Lambda_{QCD}$, $Q_0 \rightsquigarrow$ the momentum distribution of the \bar{d} and \bar{u} sea in the proton should be similar
- $\rightsquigarrow pQCD \Rightarrow x\bar{d} - x\bar{u} = 0$
- Hadronic fluctuations: (cheapest) $|n\pi^+\rangle$
- $\pi^+ \sim u\bar{d} \Rightarrow |n\pi^+\rangle \rightsquigarrow$ asymmetry
- Including $|\Delta\pi\rangle \rightsquigarrow$ even better agreement

The proton spin puzzle / “crisis”

1988 European Muon Collaboration:
measured polarization of quarks adds to only $\sim 1/3$ of proton spin $1/2$

Unpolarised:

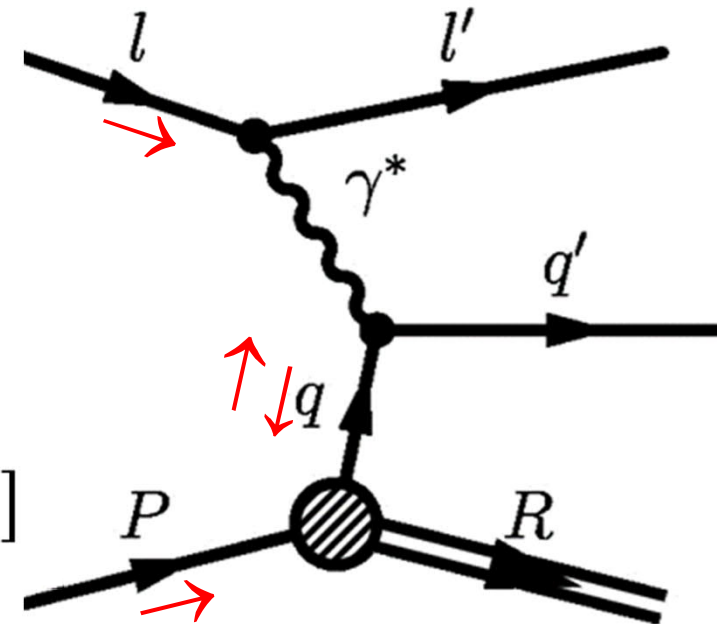
$$\frac{d^2\sigma}{dx dQ^2} \sim F_2(x, Q^2) = \sum_f e_f^2 x [q_f(x, Q^2) + \bar{q}_f(x, Q^2)]$$

$$\text{PDF: } q(x) \equiv q^\uparrow(x) + q^\downarrow(x)$$

Polarised:

$$\frac{d^2\Delta\sigma}{dx dQ^2} \sim g_1(x, Q^2) = \sum_f e_f^2 [\Delta q_f(x, Q^2) + \Delta \bar{q}_f(x, Q^2)]$$

$$\text{Polarised PDF: } \Delta q(x) \equiv q^\uparrow(x) - q^\downarrow(x)$$

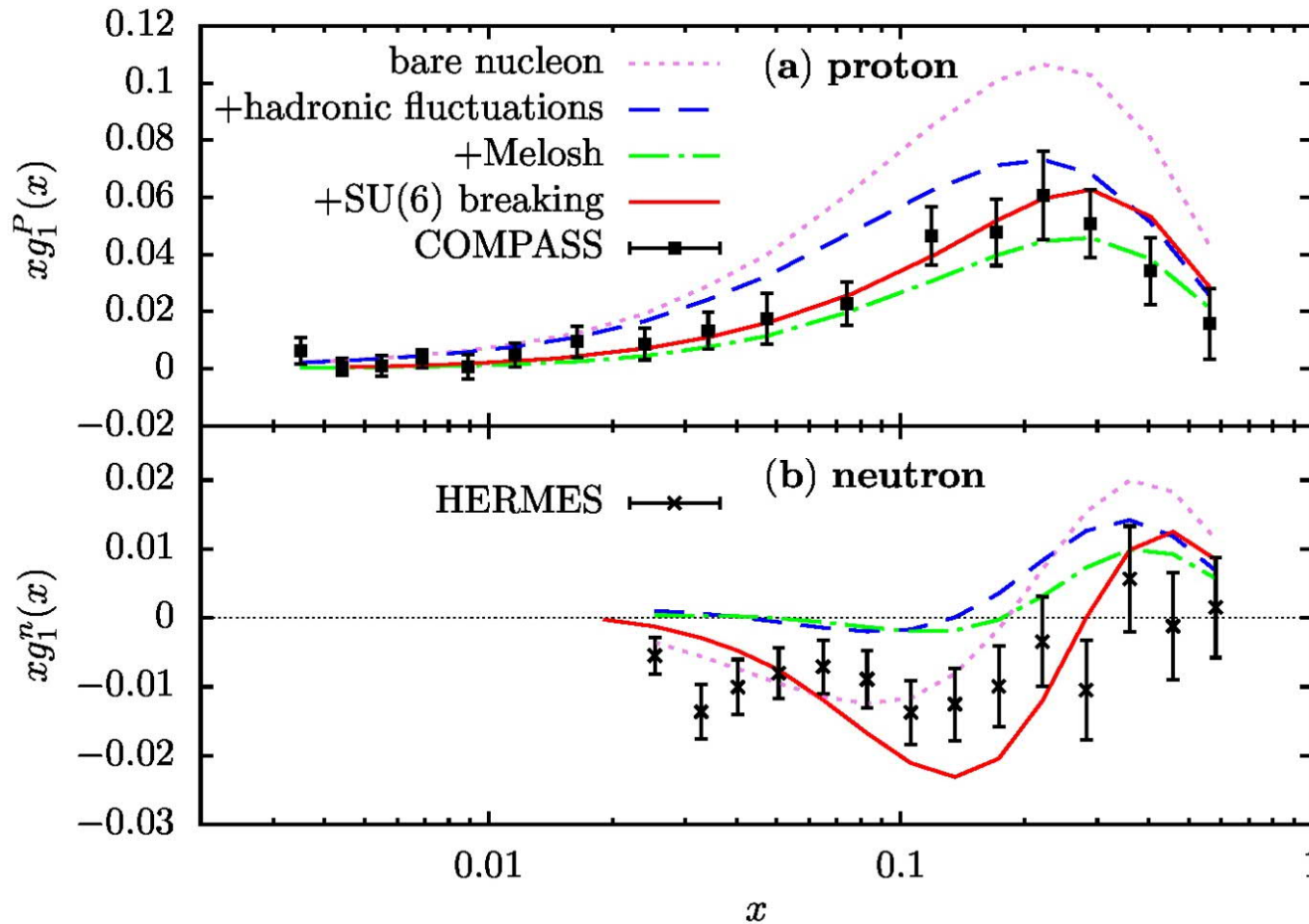


$$\rightarrow \Delta\Sigma = \int_0^1 dx \sum_f \Delta q_f(x) \approx 0.3 \text{ !?}$$

$$\text{proton spin } \frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta g + L_q + L_g$$

Polarised PDFs: $\Delta f_{q/H}^{\text{bare}}(x) = \Delta f_{q/H}^{\text{SU}(6)} f_{q/H}^{\text{bare}}(x)$

$$|p \uparrow\rangle = \frac{1}{\sqrt{2}}|u \uparrow (ud)_{S=0}\rangle + \frac{1}{\sqrt{18}}|u \uparrow (ud)_{S=1}\rangle - \frac{1}{3}|u \downarrow (ud)_{S=1}\rangle - \frac{1}{3}|d \uparrow (uu)_{S=1}\rangle + \frac{\sqrt{2}}{3}|d \downarrow (uu)_{S=1}\rangle$$



Melosh = relativistic spin transformation

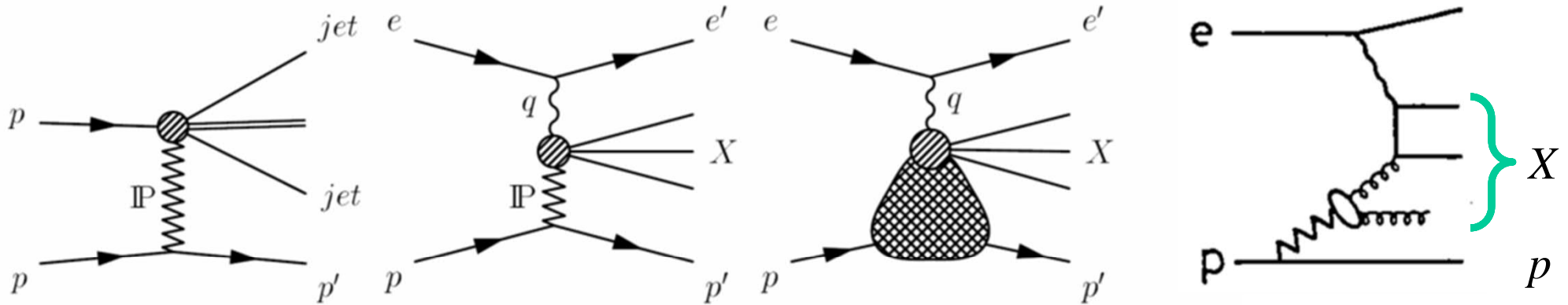
***p* spin OK !
But not *n* spin ?
Solutions ... ?**

- SU(6) breaking available range $-2 \leq \Delta f_{u/p} \leq 2$
 $-1 \leq \Delta f_{d/p} \leq 1$ $\Rightarrow \Delta f_{u/P} = 2, \Delta f_{d/P} = -1 \rightarrow |p\rangle \approx |u^\uparrow u^\uparrow d^\downarrow\rangle$
in contrast to SU(6) quark model !!??

- Problem in extracting neutron state from deuterium state ??

Possible effects of soft gluon fields missing !!

Information available from diffractive scattering



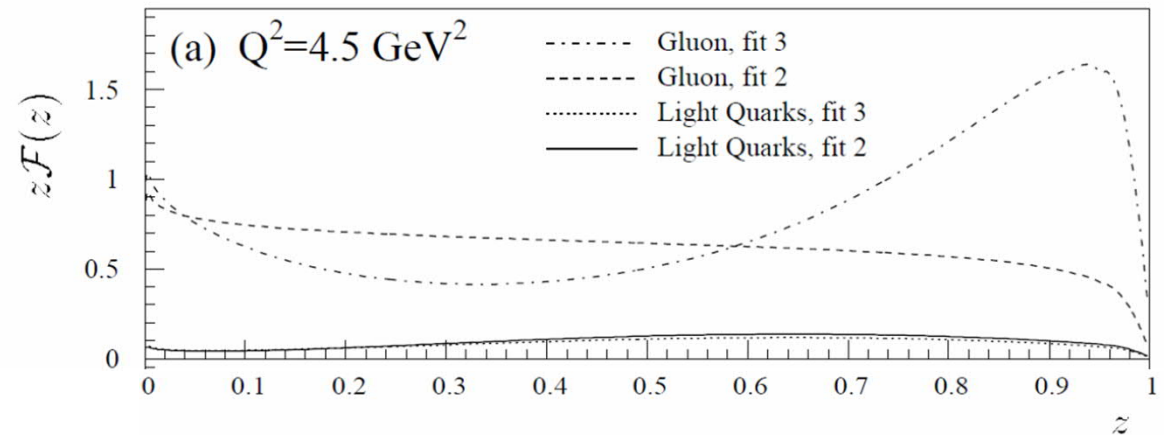
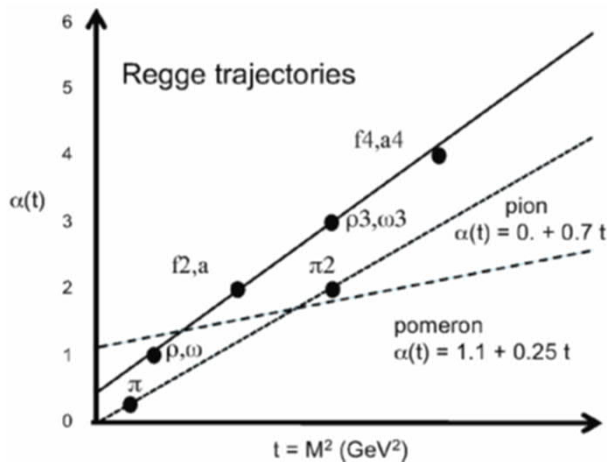
Regge theory (soft pre-QCD):

pomeron factorization $d\sigma(a + p \rightarrow p + X) = f_{IP/p}(x_{IP}, t) d\sigma(a + IP \rightarrow X)$

pomeron flux $f_{IP/p}(x_{IP}, t) = \frac{9\beta_0^2}{4\pi^2} \left(\frac{1}{x_{IP}}\right)^{2\alpha_{IP}(t)-1} [F_1(t)]^2 \quad a = p, e, \gamma \dots$

pomeron trajectory $\alpha_{IP}(t) \simeq 1.08 + 0.25t$

Pomeron as color singlet gluon field



Extra material

More info in spin sum rules:

$$\Gamma^{P\pm n}(x_{\min}) = \int_{x_{\min}}^1 dx \left(g_1^P(x) \pm g_1^n(x) \right)$$

$$\Gamma_{Bj}^{p-n}(x_{\min} = 0) = 0.187 \rightarrow$$

i.e. Bjorken sum rule

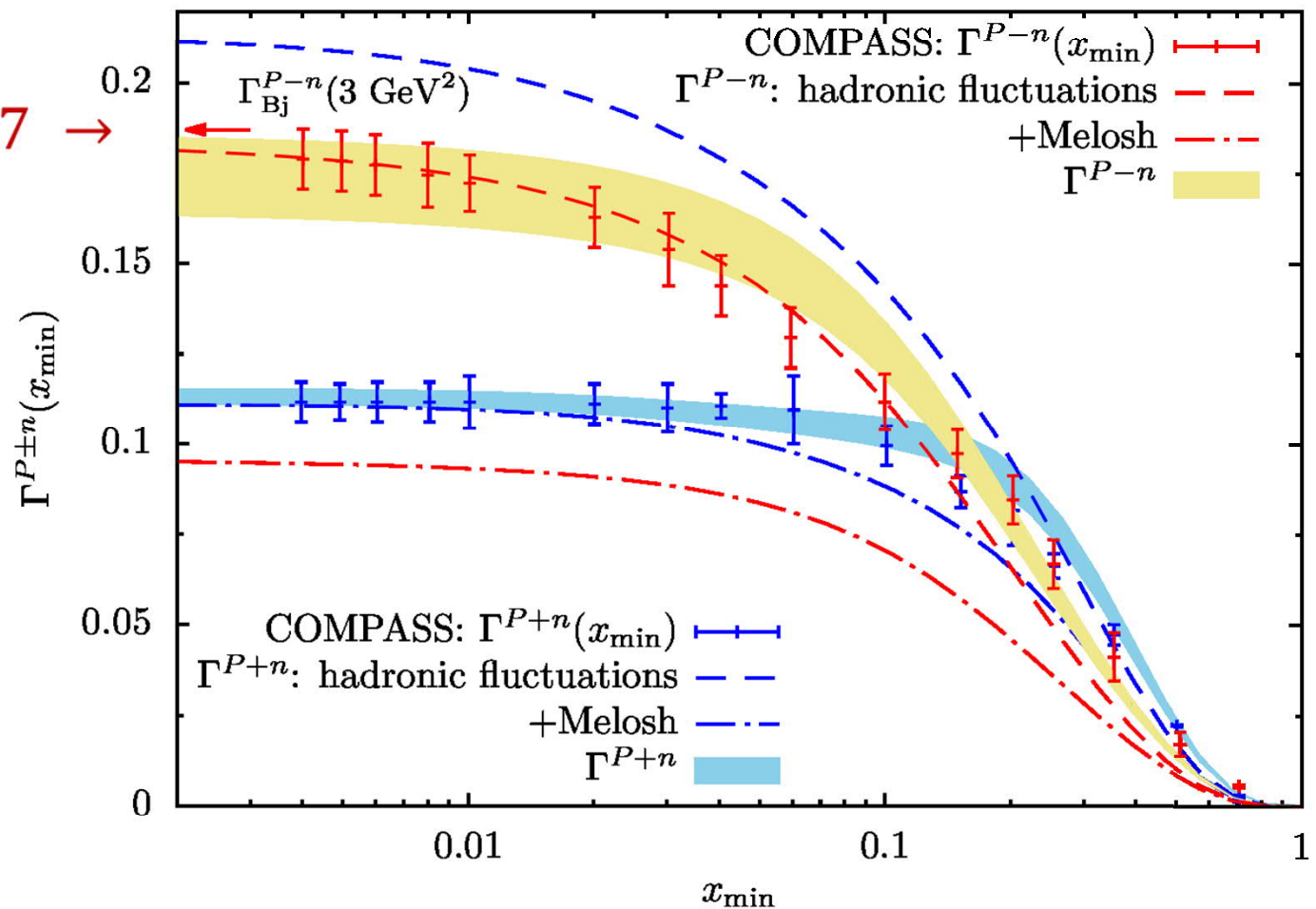
& $g_1^n(x)$ forces

SU(6) breaking,

available range:

$$-2 \leq \Delta f_{u/p} \leq 2$$

$$-1 \leq \Delta f_{d/p} \leq 1$$



$$\text{Data} \rightarrow \Gamma^{P-n}(0) \Rightarrow \Delta f_{u/P} = 2, \Delta f_{d/P} = -1$$

$\rightarrow |p\rangle \approx |u^\uparrow u^\uparrow d^\downarrow\rangle$ in contrast to SU(6) quark model !!??

Conclusion

Hadron fluctuations (ChPT) + non-pQCD parton fluctuations
+ pQCD DGLAP evolution + SU(6) quark flavor-spin symmetry
→ reproduce data on

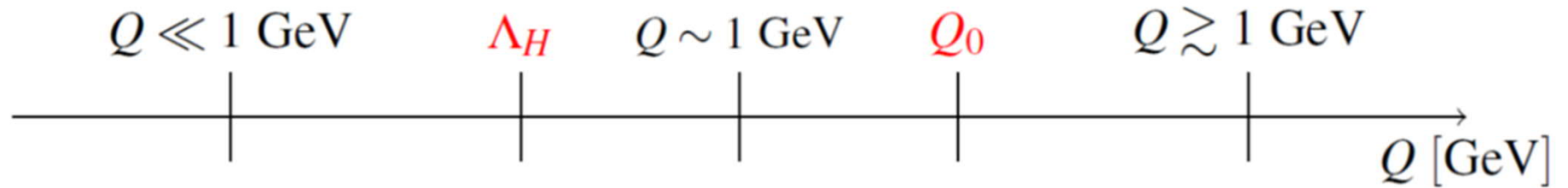
- proton spin structure function
- Ellis-Jaffe sum rule
- proton quark spin sum $\Delta\Sigma$

	bare	hadronic fluct.	Melosh	SU(6) break.	$\Delta\Sigma^{\text{exp}}$
$\Delta\Sigma$	0.95	0.75	0.39	0.39	0.26–0.36

+ SU(6) breaking → reproduce data on

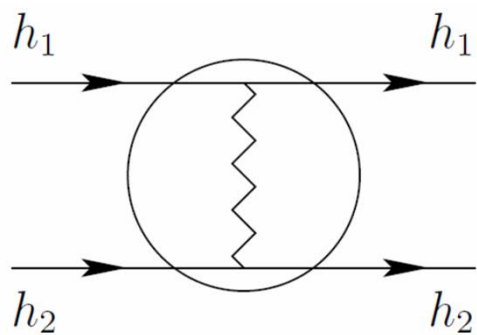
- neutron spin structure function
- Bjorken sum rule

QCD established as *the* theory for strong interactions,
 but *soft non-pQCD is the major unsolved problem* within the Standard Model

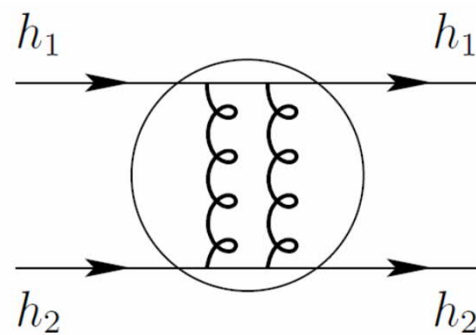


ChPT hadrons \leftarrow ??? \rightarrow partons pQCD
 Soft gluon fields !?

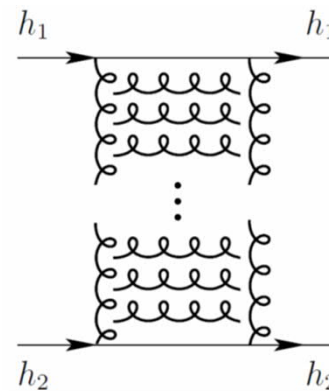
Consider information from soft elastic/diffractive hadron scattering



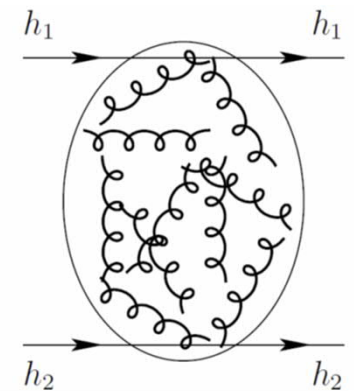
Pomeron
 theory/phenomenology



2-gluon
 exchange



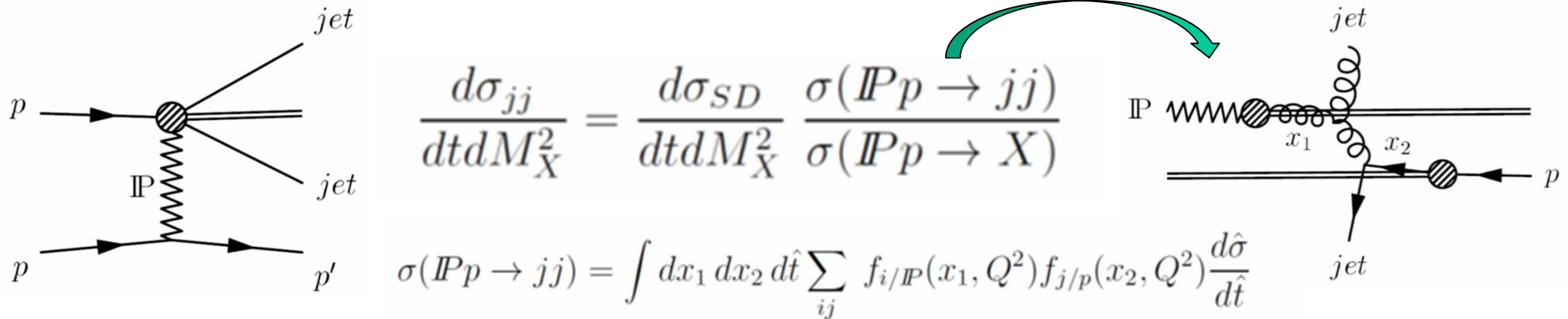
reggeized
 gluon ladder



gluon field
 vacuum fluctuation

Possible effects of soft gluon fields missing !!

Information available from diffractive hard scattering:



Based on:

pomeron factorization $d\sigma(a + p \rightarrow p + X) = f_{\mathbb{P}/p}(x_{\mathbb{P}}, t) d\sigma(a + \mathbb{P} \rightarrow X)$
 $a = p, e, \gamma \dots$

pomeron flux $f_{\mathbb{P}/p}(x_{\mathbb{P}}, t) = \frac{9\beta_0^2}{4\pi^2} \left(\frac{1}{x_{\mathbb{P}}}\right)^{2\alpha_{\mathbb{P}}(t)-1} [F_1(t)]^2$

pomeron trajectory $\alpha_{\mathbb{P}}(t) \simeq 1.08 + 0.25t$

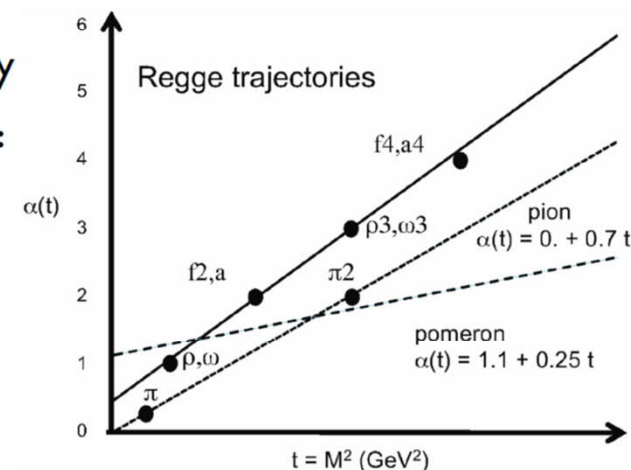
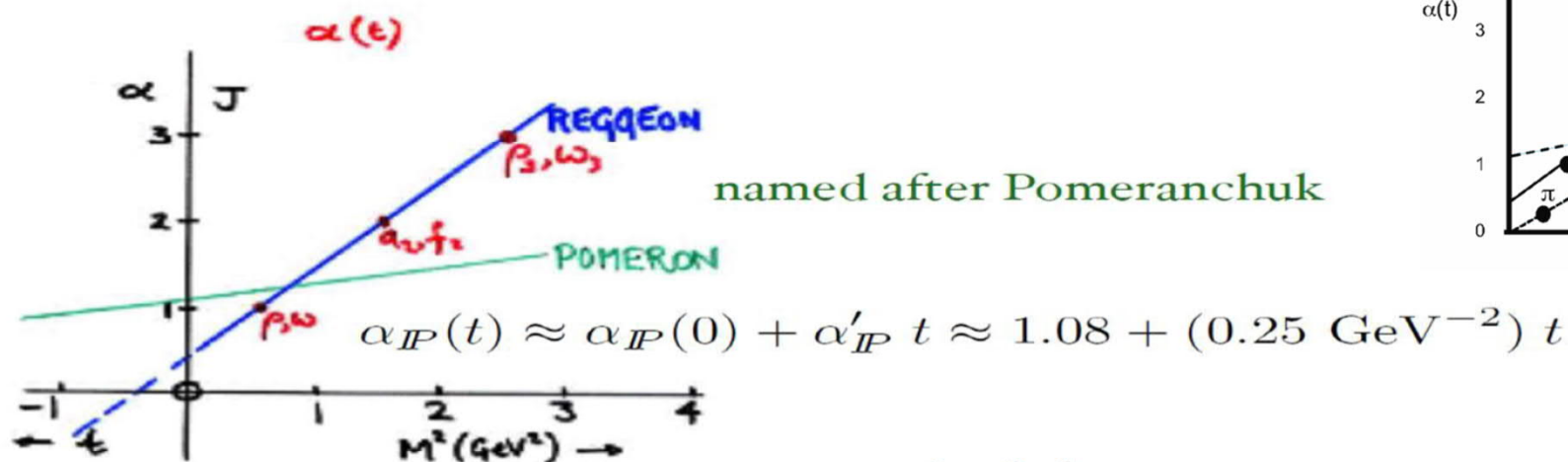
from Regge phenomenology with

$\beta = 3.24 \text{ GeV}^2$ is the mentioned effective pomeron-quark coupling and $F_1(t) = (4m_p^2 - At)/(4m_p^2 - t) \cdot (1 - t/B)^{-2}$ is a proton form factor with m_p the proton mass and parameters $A = 2.8$, $B = 0.7$.

Reggeon Field Theory before QCD

P.D.B. Collins, *An introduction to Regge theory and high energy physics*, Cambridge University Press, Cambridge, 1977.

- V.N. Gribov introduce in the 60's
- Scattering amplitude at high energies for hadrons is according Regge Theory
- The exchange are "quasi particles" characterized by its Regge trajectories :



- According to the Regge theory the contribution to the total Cross section, is given by:

$$\sigma_T = A_i S^{\alpha_i(0)-1}$$

- Leading Pole: is Called Pomeron with vacuum quantum numbers

$$\alpha(t) = \alpha_0 + \alpha' t \quad \alpha_0 \text{ is the Pomeron intercept and } \alpha' \text{ is the slope}$$

optical theorem

$$\sigma_{\text{total}} = \sum_X \left| \text{Diagram with } X \right|^2 = \text{Im} \left[\text{Diagram with } \alpha_{IP}(0) \right] = \text{Diagram with } g_N$$

Diffractive hard scattering

Idea:

Ingelman-Schlein 1984 → Blois -85!!

- hard scale probes parton level
- \mathbb{P} flux
- \mathbb{P} structure function

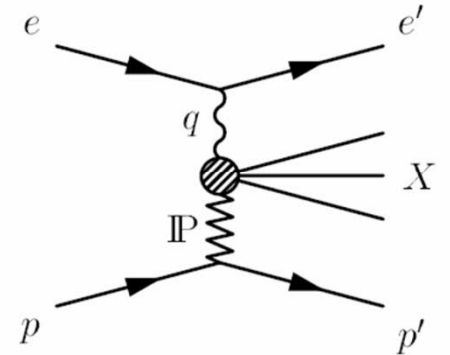
Predicted:

- jets etc. in hadronic diffraction
- diffractive deep inelastic scattering

Discovery: UA8 at $S_{p\bar{p}S}$ 1988

- jets in single diffraction \simeq model
- hard gluons $xg(x) \sim x(1-x)$ in \mathbb{P}
- ‘superhard’ $\delta(1-x)$ component

More exp's: HERA ep , Tevatron $p\bar{p}$



Diffractive DIS

$$\frac{d\sigma}{dx dQ^2 dx_{\mathbb{P}} dt} =$$

$$\frac{2\pi\alpha^2}{xQ^4} (1 + (1-y)^2) F_2^{D(4)}$$

$$F_2^{D(4)}(x, Q^2, x_{\mathbb{P}}, t) =$$

$$\underbrace{f(x_{\mathbb{P}}, t)}_{\mathbb{P} \text{ flux}} \underbrace{F_2^{\mathbb{P}}(\beta, Q^2)}_{\mathbb{P} \text{ structure}}$$

Fits HERA rapidity gap data